

Radiation and the Planck Function

Lectures in Benevento
June 2007

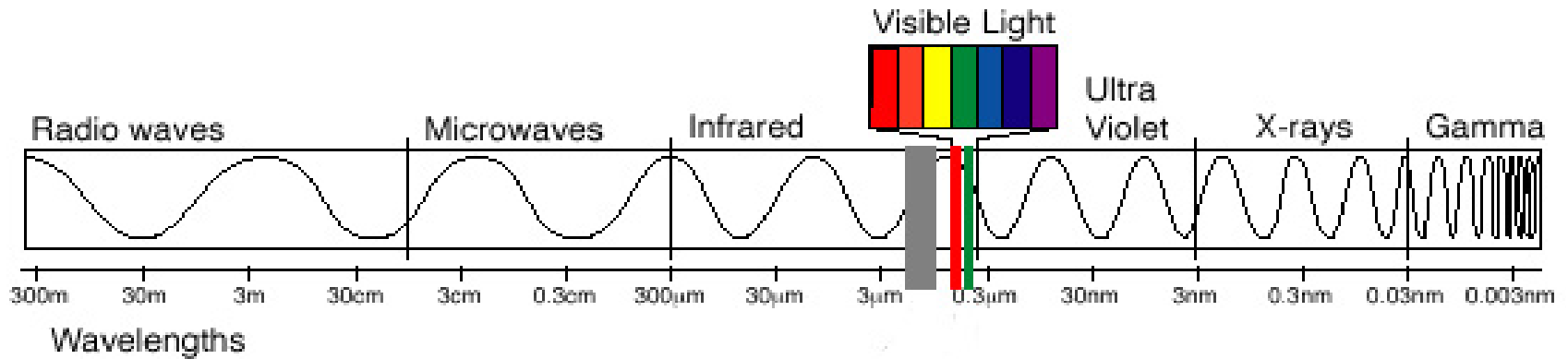
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All satellite remote sensing systems involve the measurement of electromagnetic radiation.

Electromagnetic radiation has the properties of both waves and discrete particles, although the two are never manifest simultaneously.

Electromagnetic radiation is usually quantified according to its wave-like properties; for many applications it considered to be a continuous train of sinusoidal shapes.

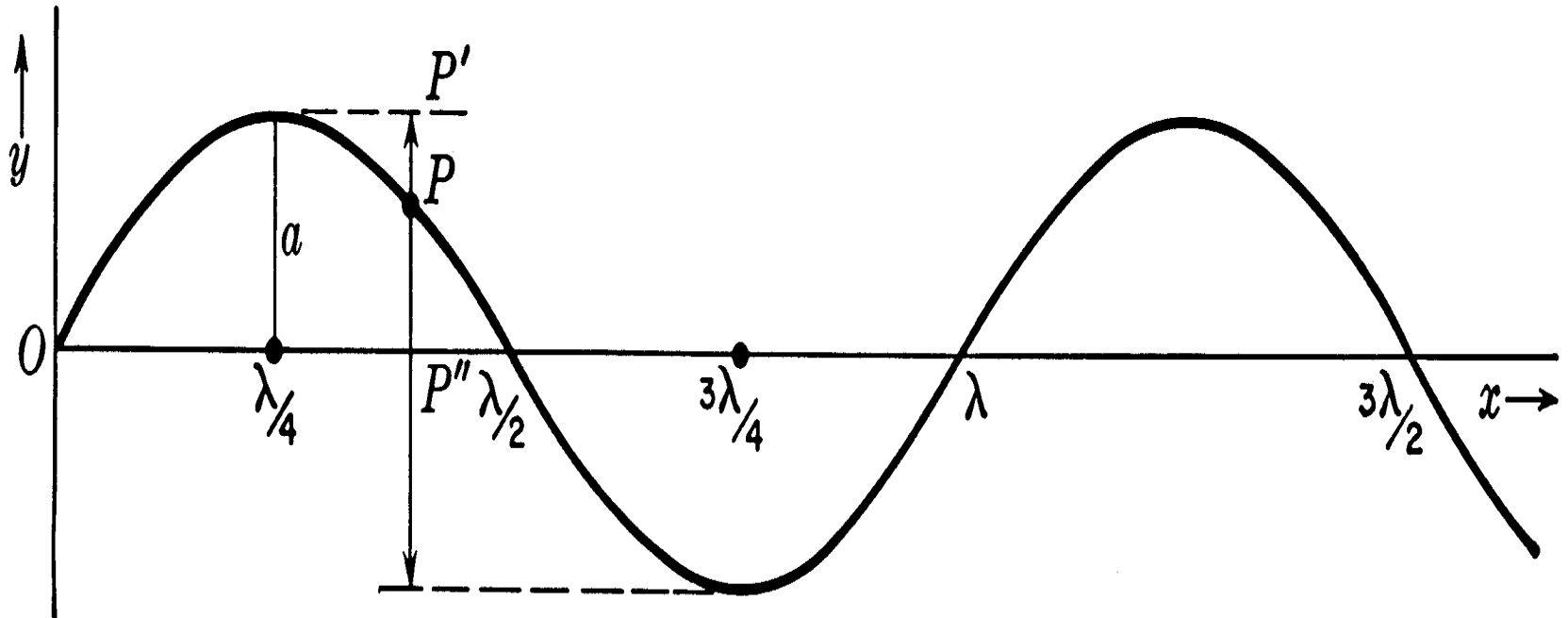
The Electromagnetic Spectrum



Remote sensing uses radiant energy that is reflected and emitted from Earth at various “wavelengths” of the electromagnetic spectrum

Our eyes are sensitive to the visible portion of the EM spectrum

Radiation is characterized by wavelength λ and amplitude a



Terminology of radiant energy

**Energy from
the Earth Atmosphere**

over time is

Flux

which strikes the detector area

Irradiance

at a given wavelength interval

**Monochromatic
Irradiance**

over a solid angle on the Earth

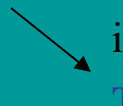
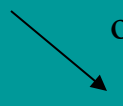
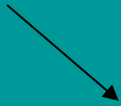
**Radiance observed by
satellite radiometer**

is described by

The Planck function

can be inverted to

Brightness temperature



Definitions of Radiation

QUANTITY	SYMBOL	UNITS
Energy	dQ	Joules
Flux	dQ/dt	Joules/sec = Watts
Irradiance	$dQ/dt/dA$	Watts/meter ²
Monochromatic Irradiance	$dQ/dt/dA/d\lambda$ or $dQ/dt/dA/d\nu$	W/m ² /micron W/m ² /cm ⁻¹
Radiance	$dQ/dt/dA/d\lambda/d\Omega$ or $dQ/dt/dA/d\nu/d\Omega$	W/m ² /micron/ster W/m ² /cm ⁻¹ /ster

Radiation from the Sun

The rate of energy transfer by electromagnetic radiation is called the radiant flux, which has units of energy per unit time. It is denoted by

$$F = dQ / dt$$

and is measured in joules per second or watts. For example, the radiant flux from the sun is about 3.90×10^{26} W.

The radiant flux per unit area is called the irradiance (or radiant flux density in some texts). It is denoted by

$$E = dQ / dt / dA$$

and is measured in watts per square metre. The irradiance of electromagnetic radiation passing through the outermost limits of the visible disk of the sun (which has an approximate radius of 7×10^8 m) is given by

$$E (\text{sun sfc}) = \frac{3.90 \times 10^{26}}{4\pi (7 \times 10^8)^2} = 6.34 \times 10^7 \text{ W m}^{-2} .$$

The solar irradiance arriving at the earth can be calculated by realizing that the flux is a constant, therefore

$$E (\text{earth sfc}) \times 4\pi R_{\text{es}}^2 = E (\text{sun sfc}) \times 4\pi R_{\text{s}}^2,$$

where R_{es} is the mean earth to sun distance (roughly 1.5×10^{11} m) and R_{s} is the solar radius. This yields

$$E (\text{earth sfc}) = 6.34 \times 10^7 (7 \times 10^8 / 1.5 \times 10^{11})^2 = 1380 \text{ W m}^{-2}.$$

The irradiance per unit wavelength interval at wavelength λ is called the monochromatic irradiance,

$$E_{\lambda} = dQ / dt / dA / d\lambda ,$$

and has the units of watts per square metre per micrometer. With this definition, the irradiance is readily seen to be

$$E = \int_0^{\infty} E_{\lambda} d\lambda .$$

In general, the irradiance upon an element of surface area may consist of contributions which come from an infinity of different directions. It is sometimes necessary to identify the part of the irradiance that is coming from directions within some specified infinitesimal arc of solid angle $d\Omega$. The irradiance per unit solid angle is called the radiance,

$$I = dQ / dt / dA / d\lambda / d\Omega,$$

and is expressed in watts per square metre per micrometer per steradian. This quantity is often also referred to as intensity and denoted by the letter B (when referring to the Planck function).

If the zenith angle, θ , is the angle between the direction of the radiation and the normal to the surface, then the component of the radiance normal to the surface is then given by $I \cos \theta$. The irradiance represents the combined effects of the normal component of the radiation coming from the whole hemisphere; that is,

$$E = \int_{\Omega} I \cos \theta d\Omega \quad \text{where in spherical coordinates } d\Omega = \sin \theta d\theta d\phi .$$

Radiation whose radiance is independent of direction is called isotropic radiation. In this case, the integration over $d\Omega$ can be readily shown to be equal to π so that

$$E = \pi I .$$

Radiation is governed by Planck's Law

$$B(\lambda, T) = \frac{c_1}{\lambda^5} \left[e^{\frac{c_2}{\lambda T}} - 1 \right]^{-1}$$

Summing the Planck function at one temperature over all wavelengths yields the energy of the radiating source

$$E = \int_{\lambda} B(\lambda, T) = \sigma T^4$$

Brightness temperature is uniquely related to radiance for a given wavelength by the Planck function.

Using wavelengths

$$\text{Planck's Law} \quad B(\lambda, T) = \frac{c_1}{\lambda^5} \left[e^{\frac{c_2}{\lambda T}} - 1 \right]^{-1} \quad (\text{mW/m}^2/\text{ster/cm})$$

where

λ = wavelengths in cm

T = temperature of emitting surface (deg K)

$c_1 = 1.191044 \times 10^{-5}$ (mW/m²/ster/cm⁻⁴)

$c_2 = 1.438769$ (cm deg K)

$$\text{Wien's Law} \quad \frac{dB(\lambda_{\max}, T)}{d\lambda} = 0 \text{ where } \lambda(\max) = .2897/T$$

indicates peak of Planck function curve shifts to shorter wavelengths (greater wavenumbers) with temperature increase. Note $B(\lambda_{\max}, T) \sim T^5$.

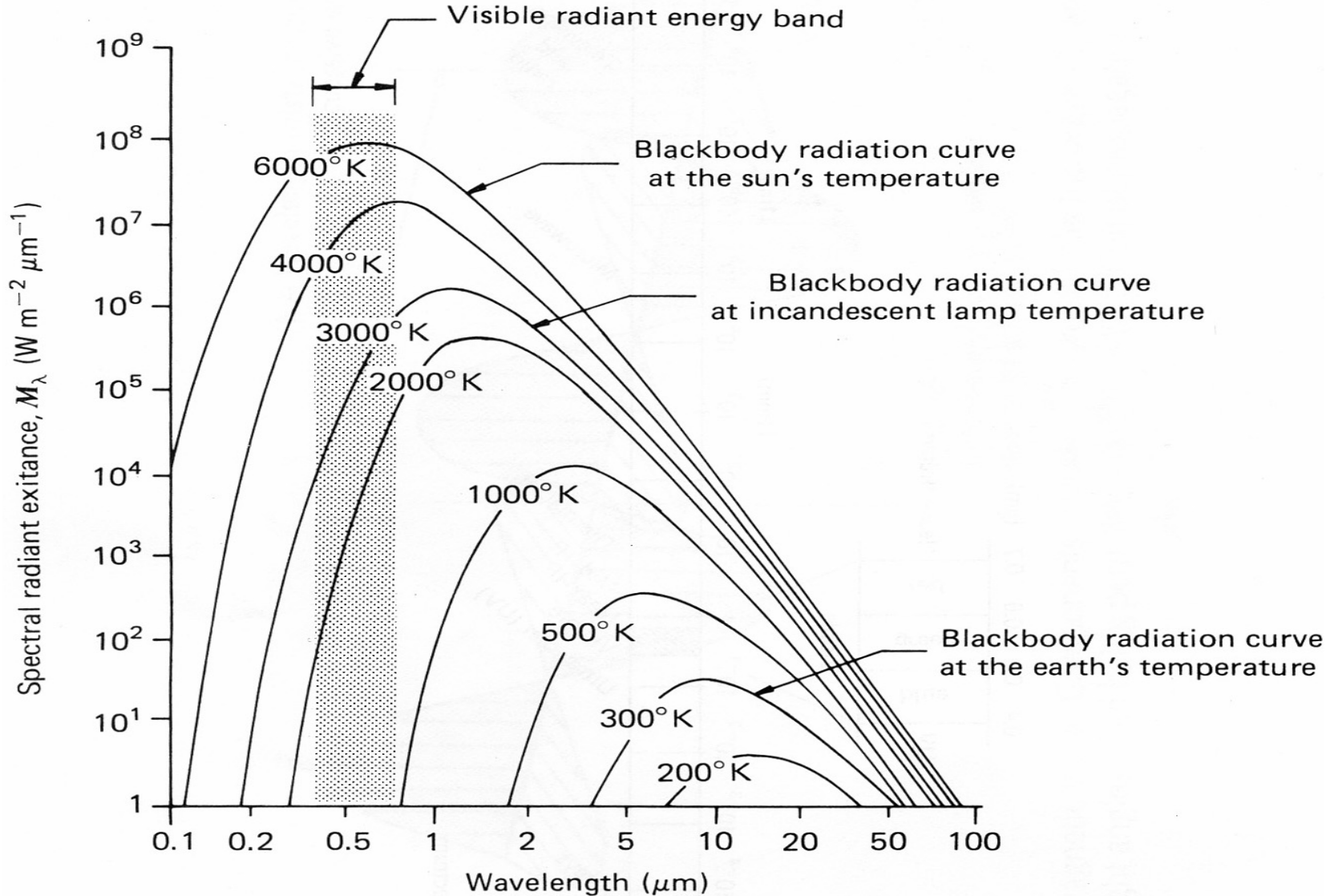
$$\text{Stefan-Boltzmann Law} \quad E = \pi \int_0^{\infty} B(\lambda, T) d\lambda = \sigma T^4, \text{ where } \sigma = 5.67 \times 10^{-8} \text{ W/m}^2/\text{deg}^4.$$

states that irradiance of a black body (area under Planck curve) is proportional to T^4 .

Brightness Temperature

$$T = \frac{c_2}{\lambda \ln\left(\frac{c_1}{\lambda^5 B_\lambda} + 1\right)} \text{ is determined by inverting Planck function}$$

Spectral Distribution of Energy Radiated from Blackbodies at Various Temperatures



Wavelength
 Wavenumber

Unnormalized
 Normalized

Wave Min
0.10

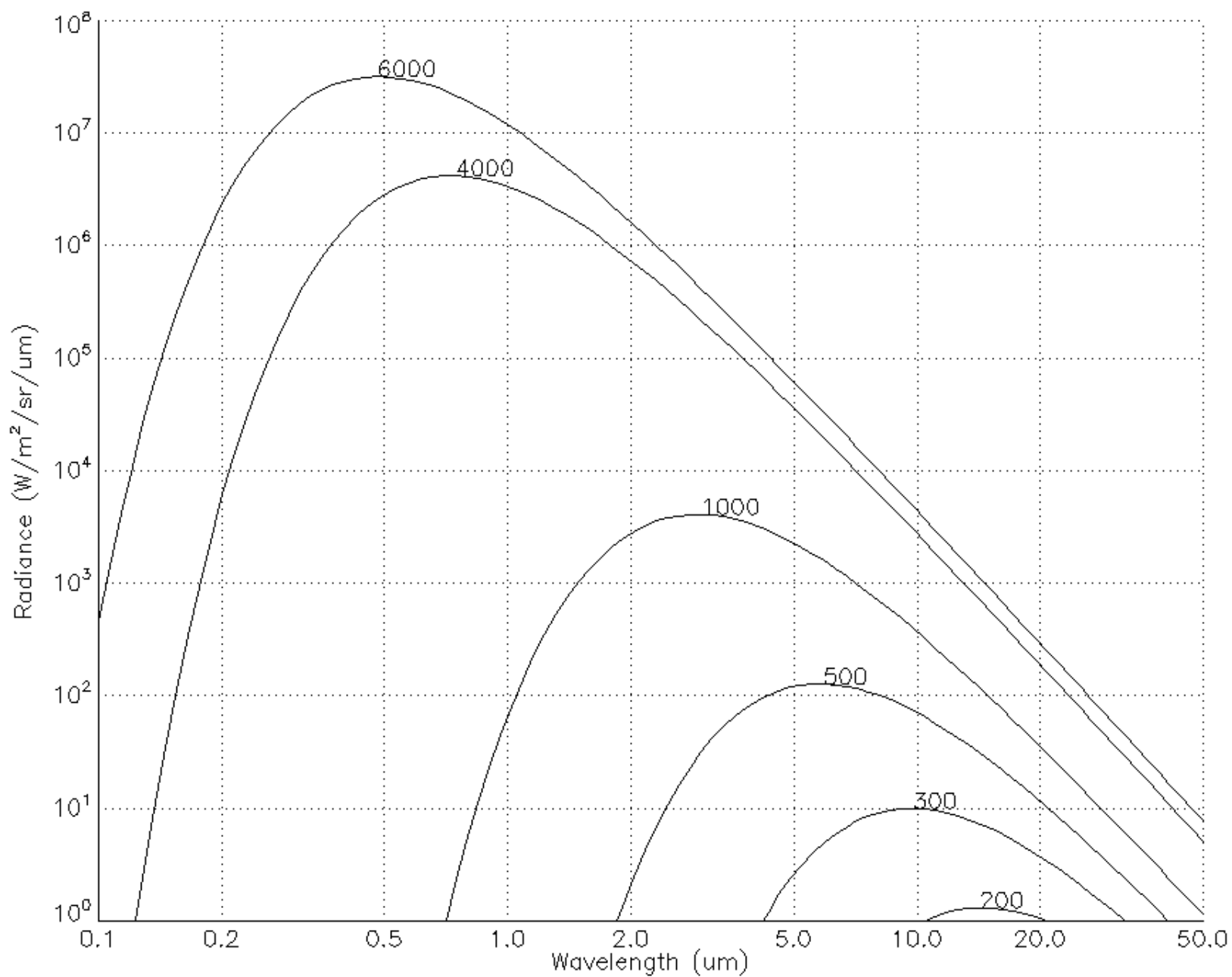
Wave Max
50.00

Temp (K)
200.00

New Plot

Add Plot

Save JPEG



$$B(\lambda_{\max}, T) \propto T^5$$

$$B(\lambda_{\max}, 6000) \sim 3.2 \times 10^7$$

$$B(\lambda_{\max}, 300) \sim 1 \times 10^1$$

so

$$B(\lambda_{\max}, 6000) / B(\lambda_{\max}, 300) \sim 3 \times 10^6$$

and

$$(6000/300)^5 = (20)^5 = 3.2 \times 10^6$$

which is the same

 Wavelength Wavenumber Unnormalized Normalized

Wave Min

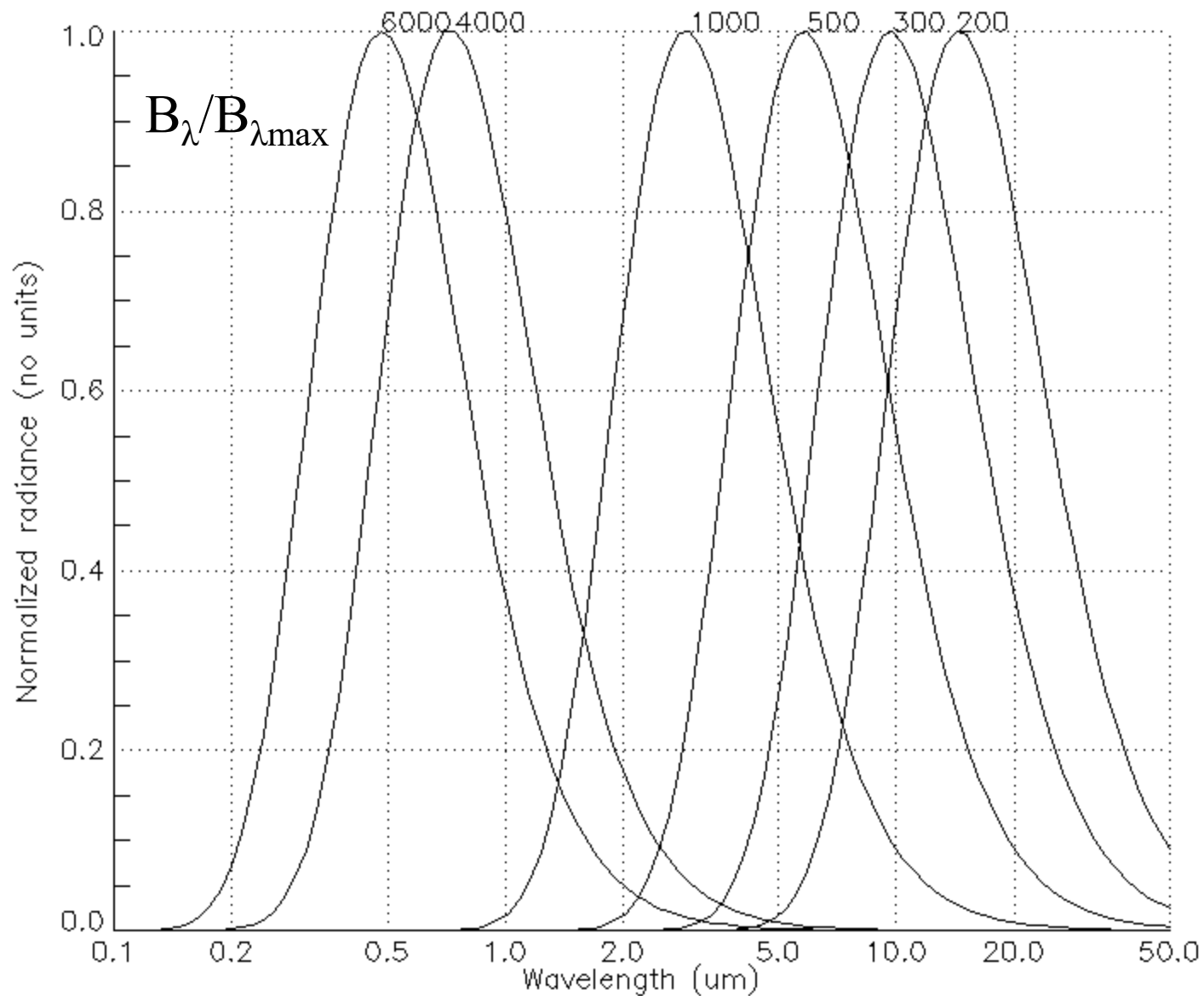
Wave Max

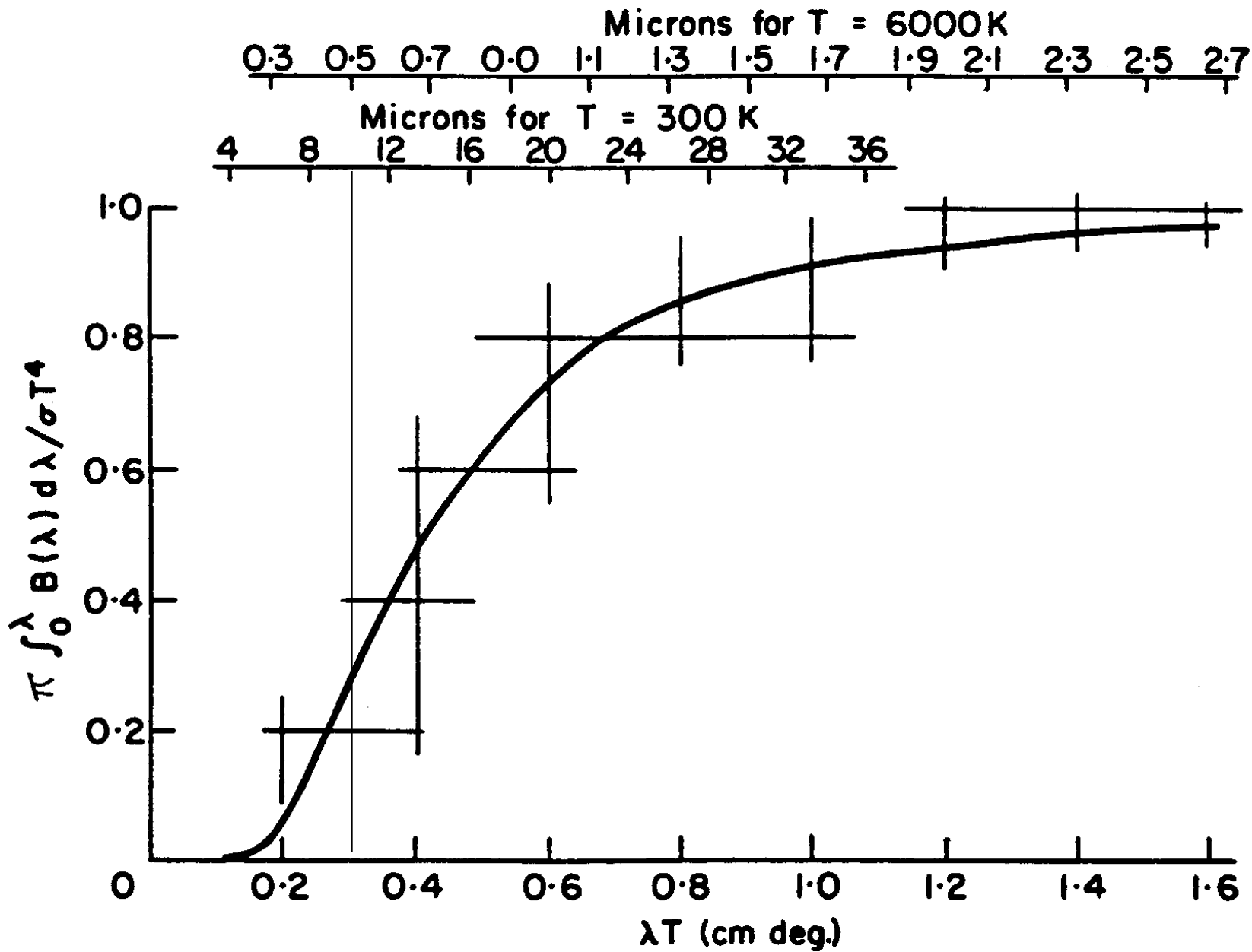
Temp (K)

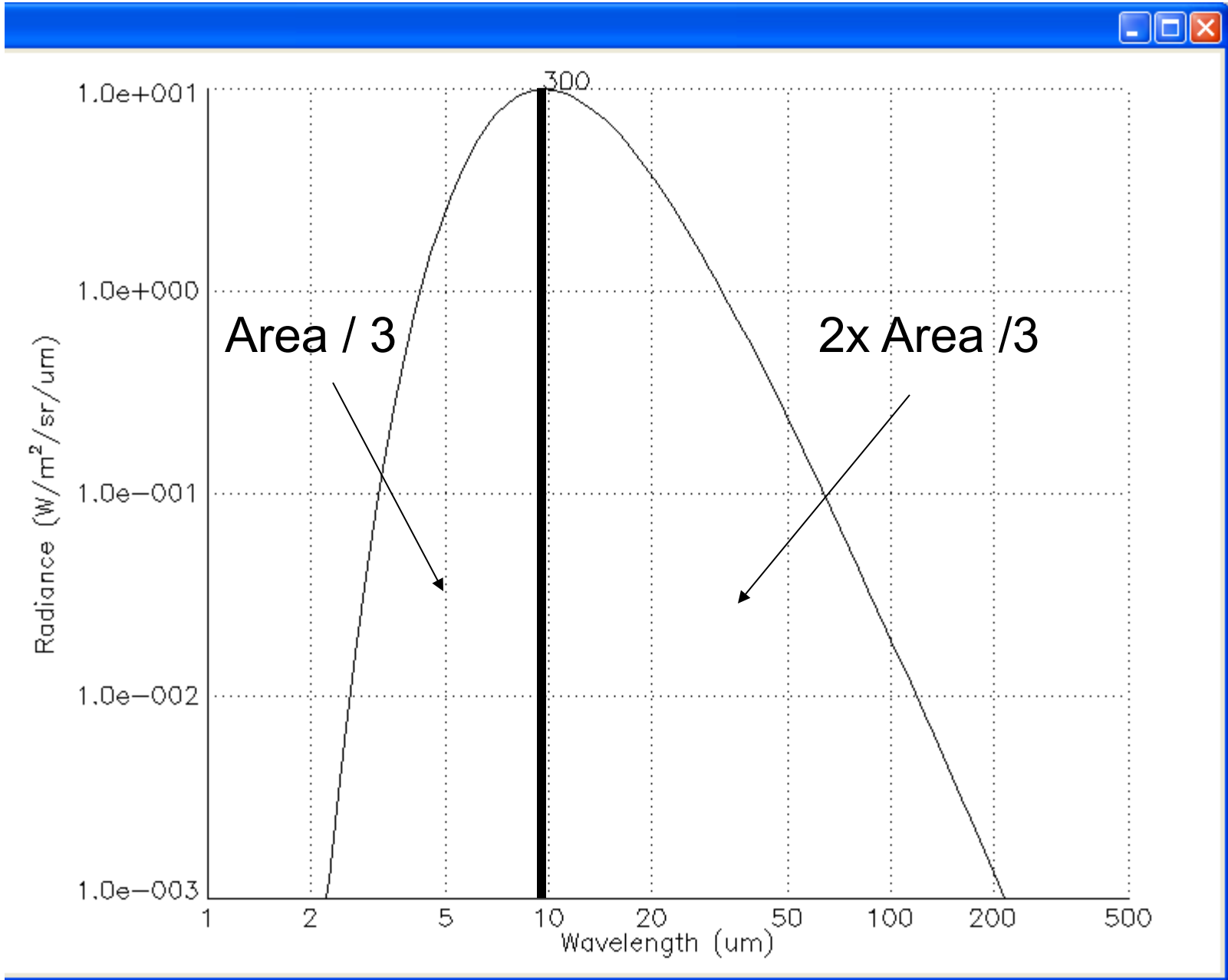
New Plot

Add Plot

Save JPEG







Using wavenumbers

Planck's Law
$$B(\nu, T) = \frac{c_1 \nu^3}{[e^{c_2 \nu / T} - 1]} \quad (\text{mW/m}^2/\text{ster/cm}^{-1})$$

where $\nu = \# \text{ wavelengths in one centimeter (cm}^{-1}\text{)}$

$T = \text{temperature of emitting surface (deg K)}$

$c_1 = 1.191044 \times 10^{-5} \text{ (mW/m}^2/\text{ster/cm}^{-4}\text{)}$

$c_2 = 1.438769 \text{ (cm deg K)}$

Wien's Law
$$dB(\nu_{\max}, T) / d\nu = 0 \text{ where } \nu(\max) = 1.95T$$

indicates peak of Planck function curve shifts to shorter wavelengths (greater wavenumbers) with temperature increase. Note $B(\nu_{\max}, T) \sim T^{**3}$.

Stefan-Boltzmann Law
$$E = \pi \int_0^{\infty} B(\nu, T) d\nu = \sigma T^4, \text{ where } \sigma = 5.67 \times 10^{-8} \text{ W/m}^2/\text{deg}^4.$$

states that irradiance of a black body (area under Planck curve) is proportional to T^4 .

Brightness Temperature

$$T = \frac{c_2 \nu}{[\ln(\frac{c_1 \nu^3}{B_\nu} + 1)]}$$

is determined by inverting Planck function

Wavelength Wavenumber Unnormalized Normalized

Wave Min

10.00

Wave Max

10000.00

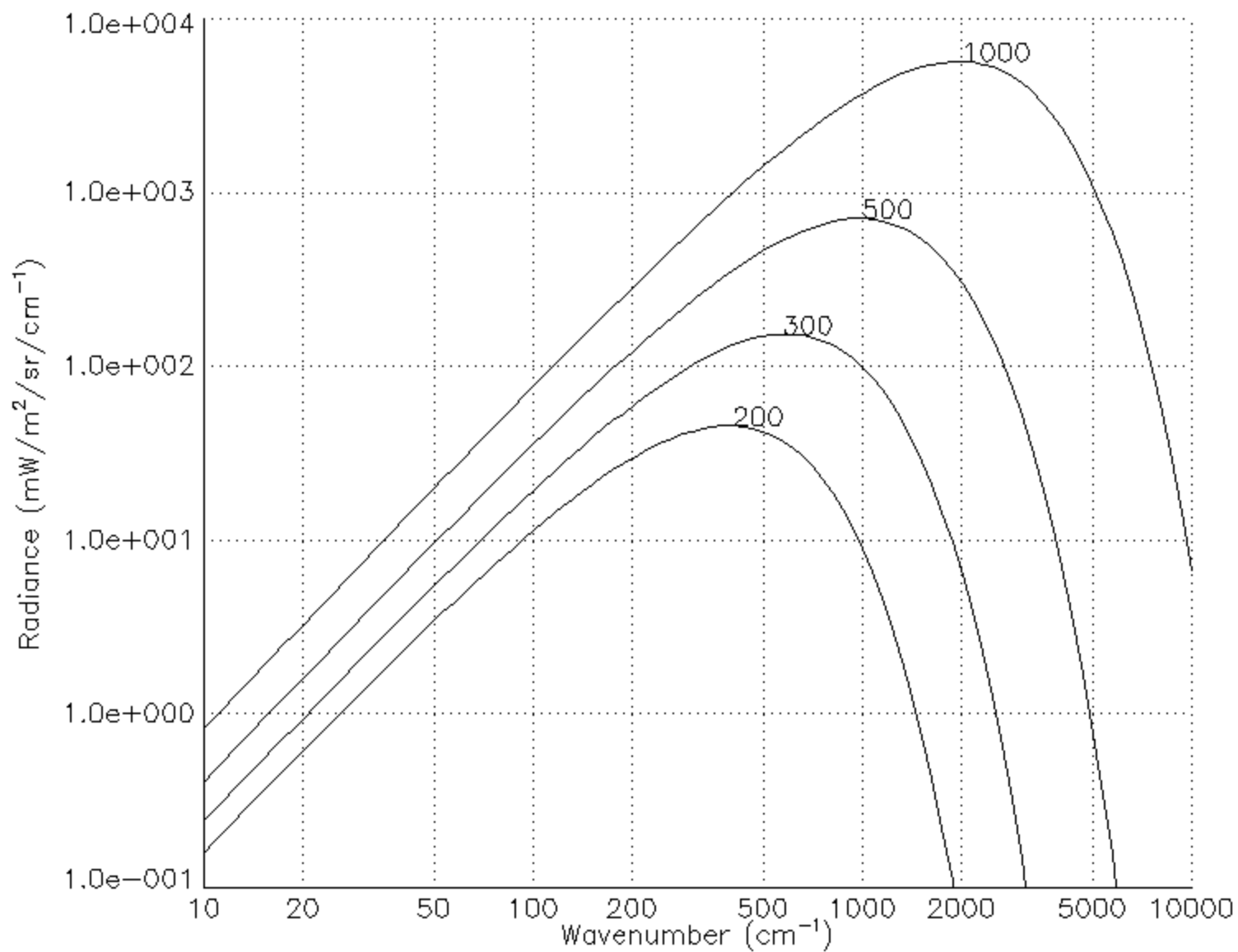
Temp (K)

200.00

New Plot

Add Plot

Save JPEG



Wavelength Wavenumber Unnormalized Normalized

Wave Min

10.00

Wave Max

10000.00

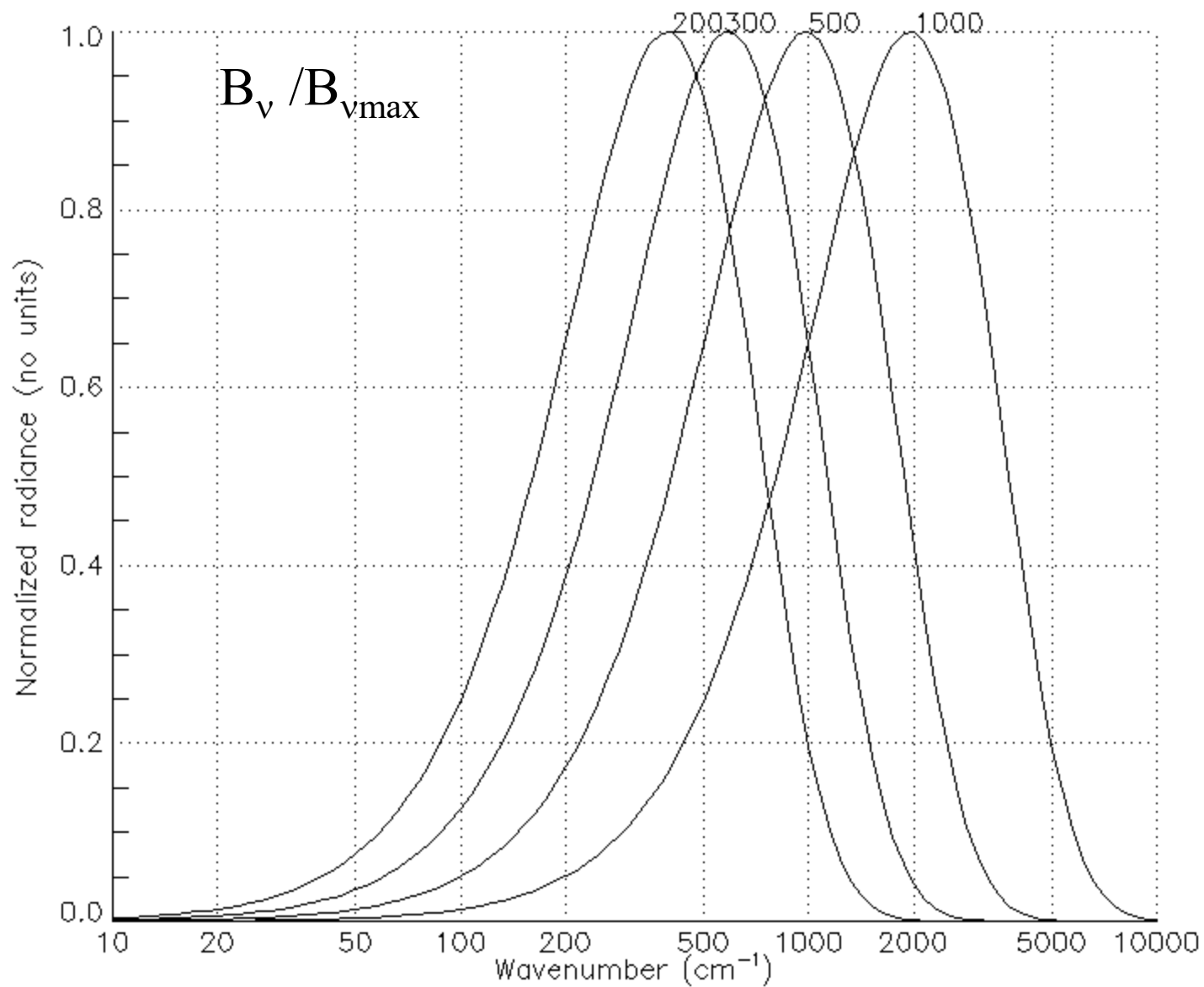
Temp (K)

200.00

New Plot

Add Plot

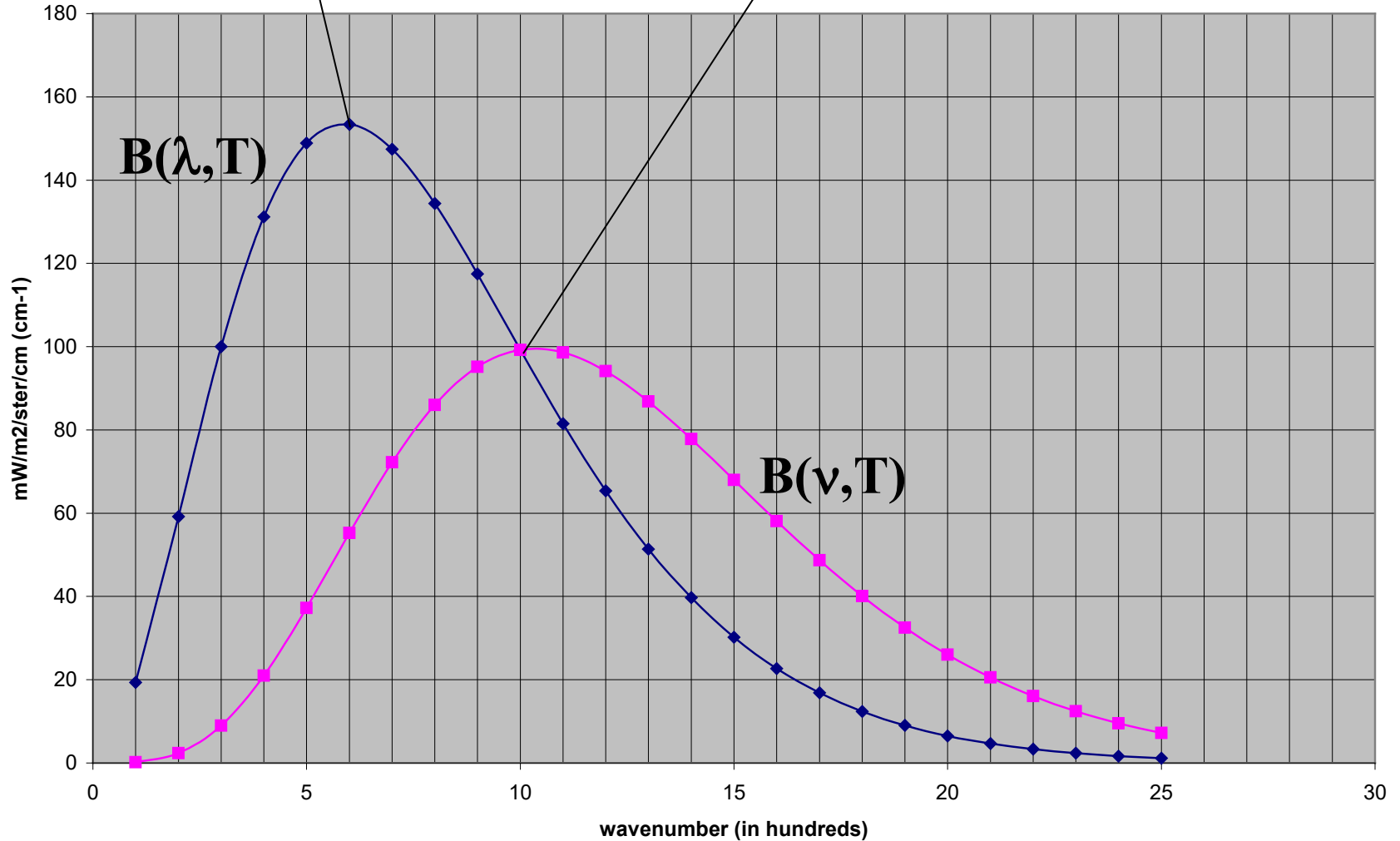
Save JPEG



$$B(\lambda_{\max}, T) \sim T^5$$

$$B(\nu_{\max}, T) \sim T^3$$

Planck Radiances



B(λ, T) versus B(ν, T)

Using wavenumbers

$$B(\nu, T) = \frac{c_1 \nu^3}{e^{c_2 \nu / T} - 1} \quad (\text{mW/m}^2/\text{ster/cm}^{-1})$$

$$\nu(\text{max in cm}^{-1}) = 1.95T$$

$$B(\nu_{\text{max}}, T) \sim T^{**3}.$$

$$E = \pi \int_0^{\infty} B(\nu, T) d\nu = \sigma T^4,$$

$$T = \frac{c_2 \nu}{\left[\ln\left(\frac{c_1 \nu^3}{B_\nu} + 1\right) \right]}$$

Using wavelengths

$$B(\lambda, T) = \frac{c_1}{\lambda^5 \left[e^{c_2 / \lambda T} - 1 \right]} \quad (\text{mW/m}^2/\text{ster}/\mu\text{m})$$

$$\lambda(\text{max in cm})T = 0.2897$$

$$B(\lambda_{\text{max}}, T) \sim T^{**5}.$$

$$E = \pi \int_0^{\infty} B(\lambda, T) d\lambda = \sigma T^4,$$

$$T = \frac{c_2}{\left[\lambda \ln\left(\frac{c_1}{\lambda^5 B_\lambda} + 1\right) \right]}$$

Temperature sensitivity, or the percentage change in radiance corresponding to a percentage change in temperature, α , is defined as

$$dB/B = \alpha dT/T.$$

The temperature sensitivity indicates the power to which the Planck radiance depends on temperature, since B proportional to T^α satisfies the equation. For infrared wavelengths,

$$\alpha = c_2\nu/T = c_2/\lambda T.$$

Wavenumber	Typical Scene Temperature	Temperature Sensitivity
700	220	4.58
900	300	4.32
1200	300	5.76
1600	240	9.59
2300	220	15.04
2500	300	11.99

$\frac{dB}{B} = \alpha \frac{dT}{T}$ or $B = c T^\alpha$ where $\alpha = \frac{c_2}{\lambda T}$ for a small temperature window around T .

$$B = B(T_0) + \left(\frac{dB}{dT}\right)_0 (\Delta T) + \left(\frac{d^2B}{dT^2}\right)_0 (\Delta T)^2 + O(3)$$

negligible

So to first order

$$c (T_0 + \Delta T)^\alpha = c T_0^\alpha + c \alpha T_0^{\alpha-1} (\Delta T)$$

$$c (T_0 + \Delta T)^\alpha - c T_0^\alpha = c \alpha T_0^{\alpha-1} (\Delta T)$$

$$\Delta B = c \alpha T_0^{\alpha-1} (\Delta T)$$

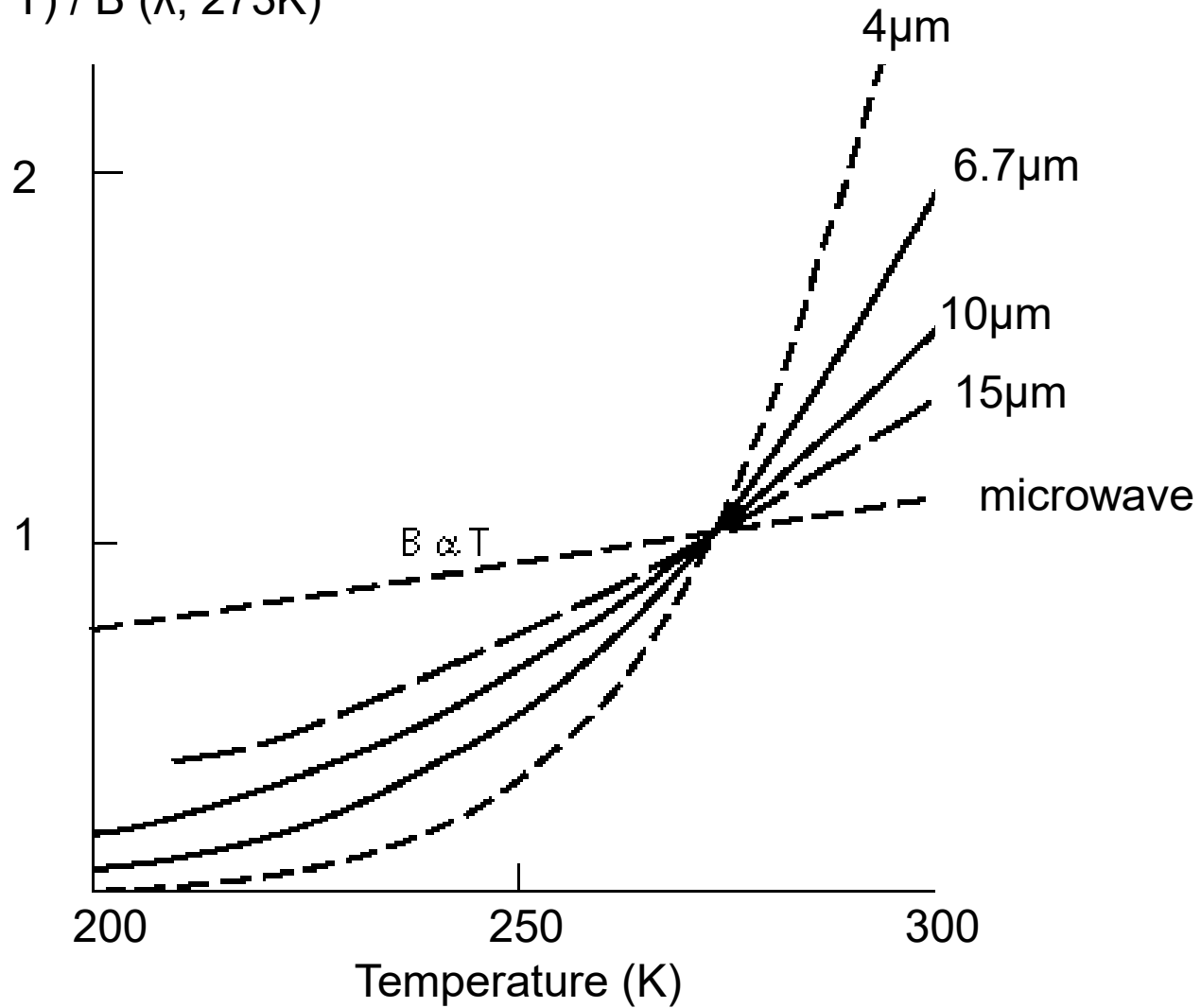
$$\frac{\Delta B}{B} = \alpha \frac{\Delta T}{T}$$

Also to first order

$$(T_0 + \Delta T)^\alpha = T_0^\alpha + \alpha T_0^{\alpha-1} (\Delta T)$$

Temperature Sensitivity of $B(\lambda, T)$ for typical earth scene temperatures

$B(\lambda, T) / B(\lambda, 273K)$



$$B(10 \text{ } \mu\text{m}, T) / B(10 \text{ } \mu\text{m}, 273) \propto T^4$$

$$B(10 \text{ } \mu\text{m}, 273) = 6.1$$

$$B(10 \text{ } \mu\text{m}, 200) = 0.9 \rightarrow 0.15$$

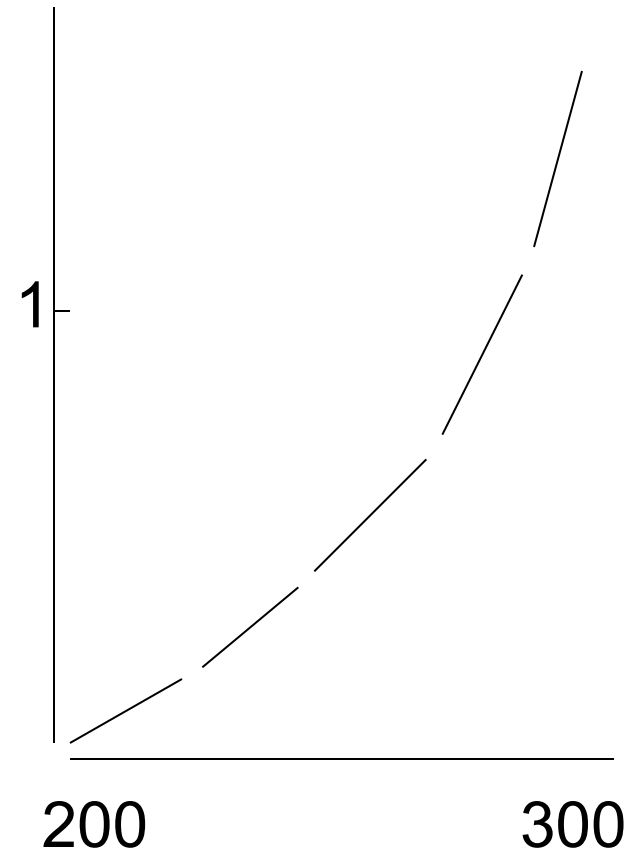
$$B(10 \text{ } \mu\text{m}, 220) = 1.7 \rightarrow 0.28$$

$$B(10 \text{ } \mu\text{m}, 240) = 3.0 \rightarrow 0.49$$

$$B(10 \text{ } \mu\text{m}, 260) = 4.7 \rightarrow 0.77$$

$$B(10 \text{ } \mu\text{m}, 280) = 7.0 \rightarrow 1.15$$

$$B(10 \text{ } \mu\text{m}, 273) = 9.9 \rightarrow 1.62$$



$$B(4 \text{ um}, T) / B(4 \text{ um}, 273) \propto T^{12}$$

$$B(4 \text{ um}, 273) = 2.2 \times 10^{-1}$$

$$B(4 \text{ um}, 200) = 1.8 \times 10^{-3} \rightarrow 0.0$$

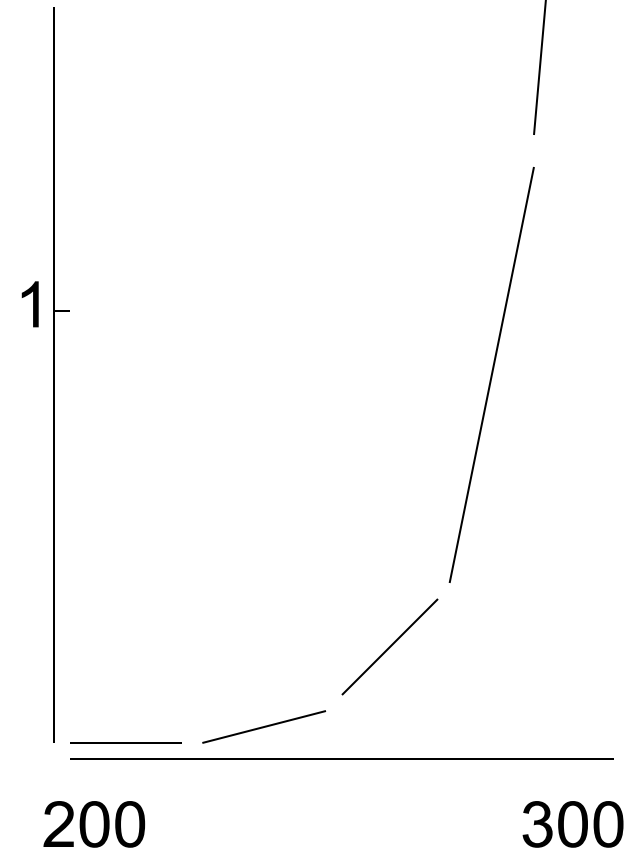
$$B(4 \text{ um}, 220) = 9.2 \times 10^{-3} \rightarrow 0.0$$

$$B(4 \text{ um}, 240) = 3.6 \times 10^{-2} \rightarrow 0.2$$

$$B(4 \text{ um}, 260) = 1.1 \times 10^{-1} \rightarrow 0.5$$

$$B(4 \text{ um}, 280) = 3.0 \times 10^{-1} \rightarrow 1.4$$

$$B(4 \text{ um}, 273) = 7.2 \times 10^{-1} \rightarrow 3.3$$



$$B(0.3 \text{ cm}, T) / B(0.3 \text{ cm}, 273) \propto T$$

$$B(0.3 \text{ cm}, 273) = 2.55 \times 10^{-4}$$

$$B(0.3 \text{ cm}, 200) = 1.8 \rightarrow 0.7$$

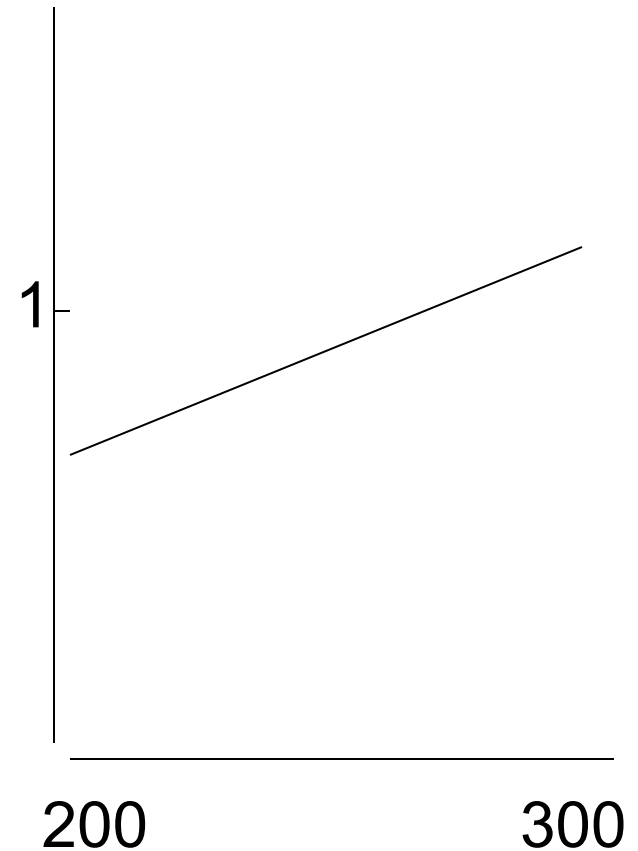
$$B(0.3 \text{ cm}, 220) = 2.0 \rightarrow 0.78$$

$$B(0.3 \text{ cm}, 240) = 2.2 \rightarrow 0.86$$

$$B(0.3 \text{ cm}, 260) = 2.4 \rightarrow 0.94$$

$$B(0.3 \text{ cm}, 280) = 2.6 \rightarrow 1.02$$

$$B(0.3 \text{ cm}, 273) = 2.8 \rightarrow 1.1$$



Radiation is governed by Planck's Law

$$B(\lambda, T) = \frac{c_1}{\lambda^5} \left[e^{\frac{c_2}{\lambda T}} - 1 \right]^{-1}$$

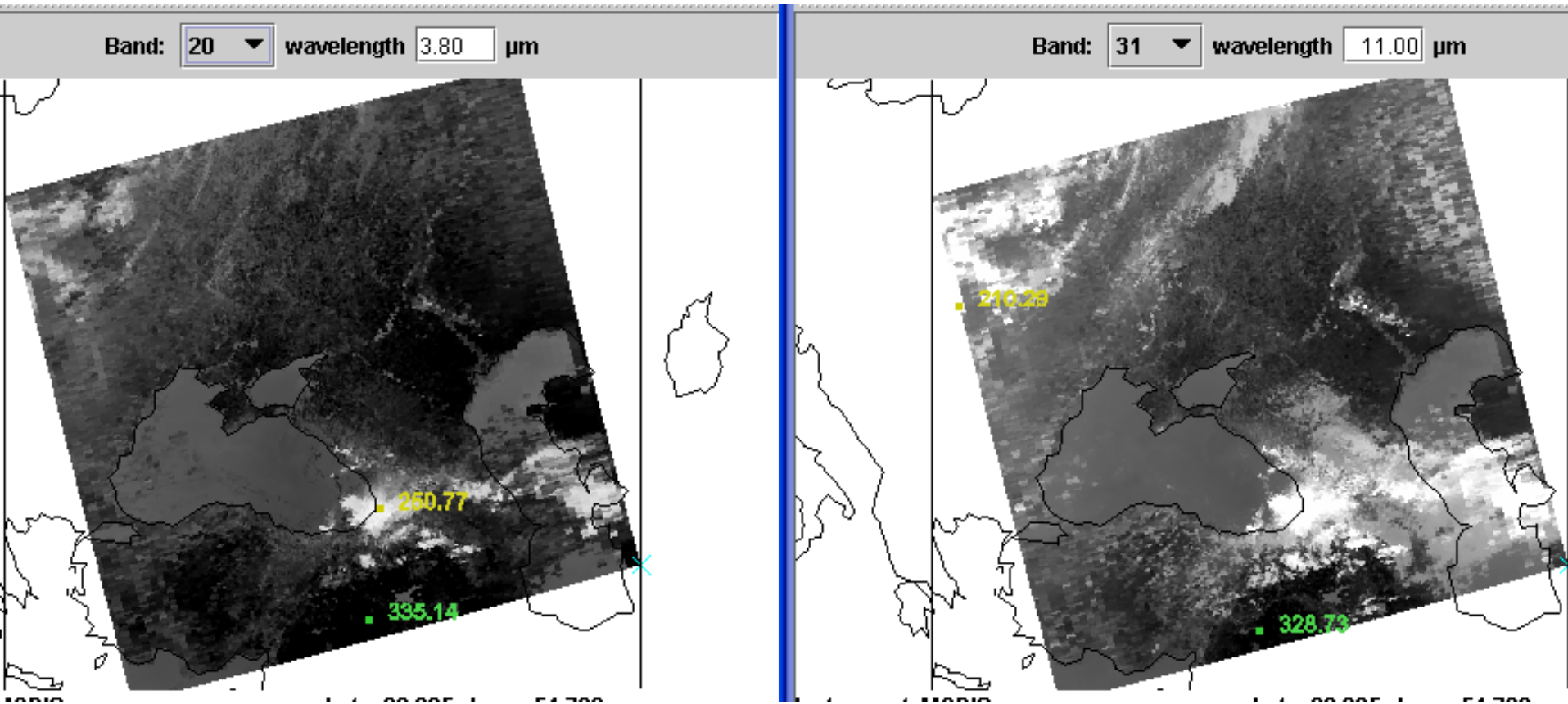
In microwave region $c_2/\lambda T \ll 1$ so that

$$e^{\frac{c_2}{\lambda T}} = 1 + \frac{c_2}{\lambda T} + \text{second order}$$

And classical Rayleigh Jeans radiation equation emerges

$$B_\lambda(T) \approx \left[\frac{c_1}{c_2} \right] \left[\frac{T}{\lambda^4} \right]$$

Radiance is linear function of brightness temperature.



Cloud edges and broken clouds appear different in 11 and 4 um images.

$$T(11)^{**4} = (1-N) * T_{clr}^{**4} + N * T_{cld}^{**4} \sim (1-N) * 300^{**4} + N * 200^{**4}$$

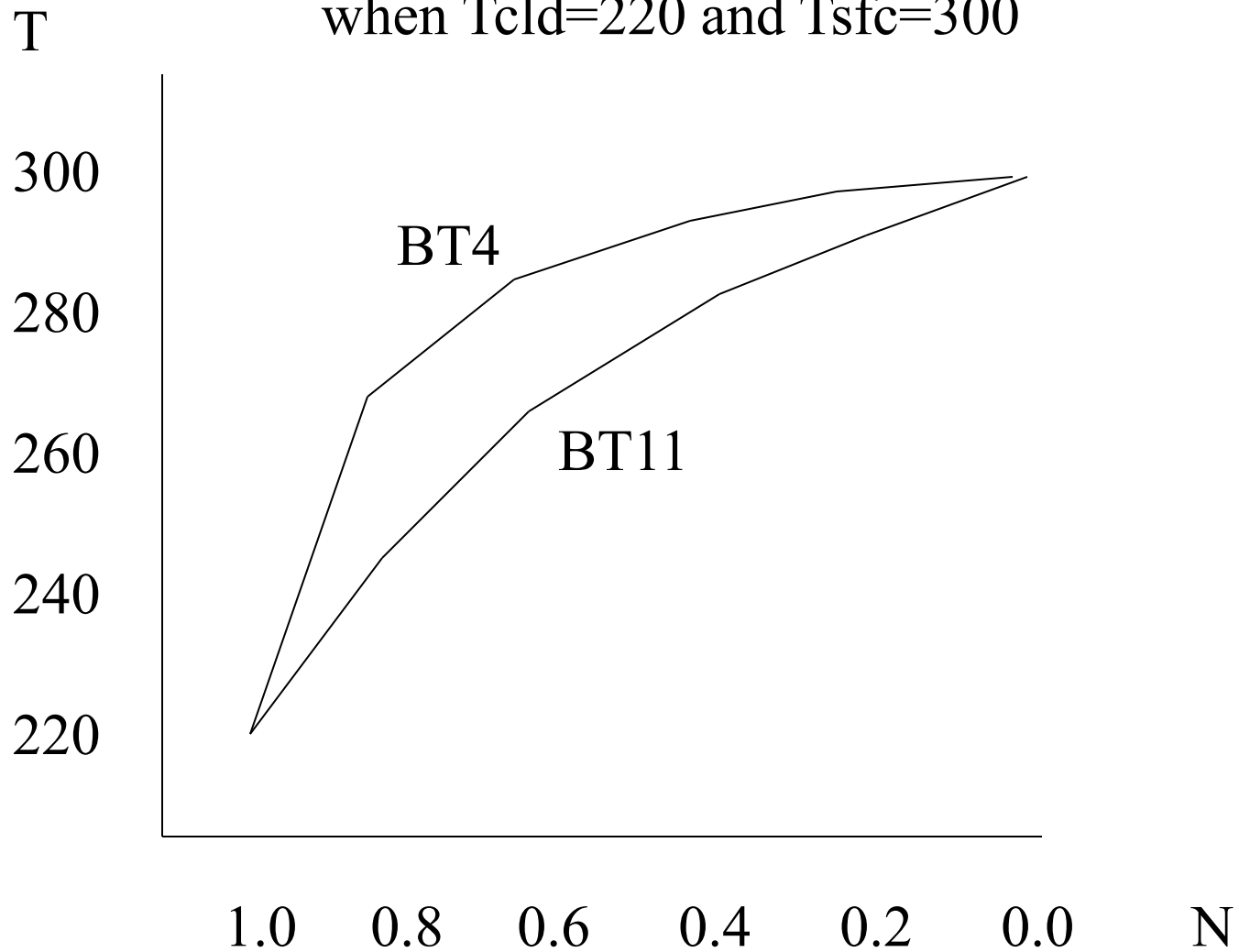
$$T(4)^{**12} = (1-N) * T_{clr}^{**12} + N * T_{cld}^{**12} \sim (1-N) * 300^{**12} + N * 200^{**12}$$

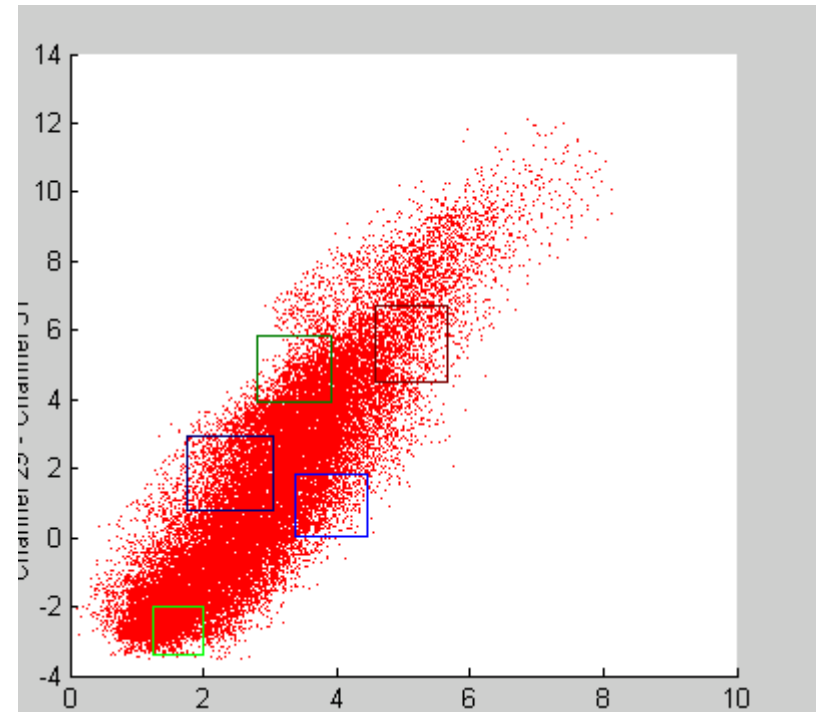
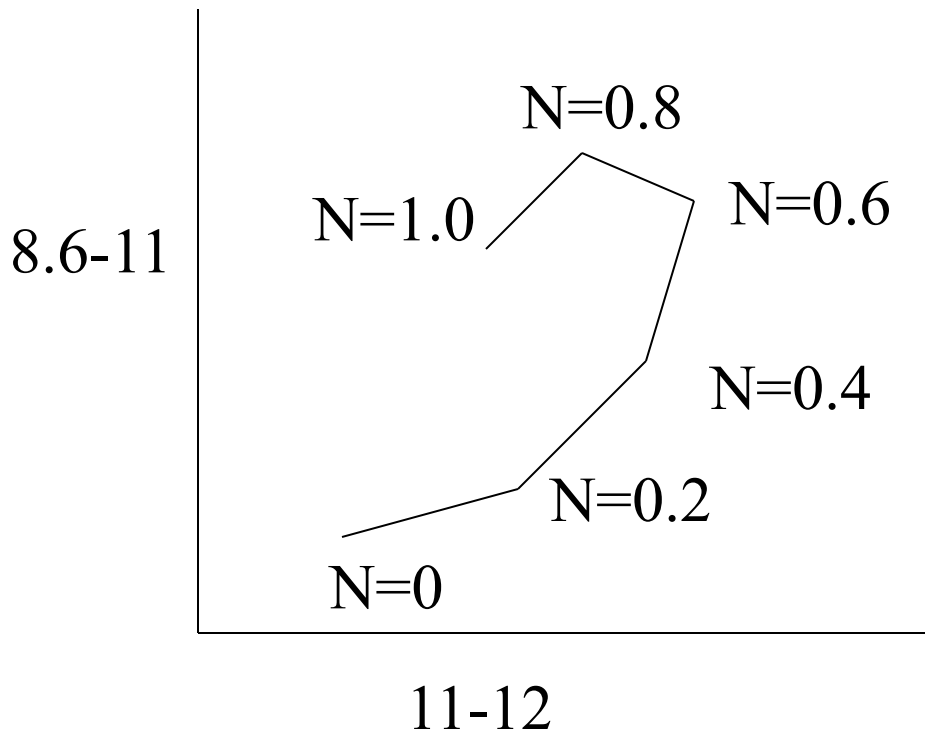
Cold part of pixel has more influence for B(11) than B(4)

Table 6.1 Longwave and Shortwave Window Planck Radiances ($\text{mW/m}^2/\text{ster/cm}^{-1}$) and Brightness Temperatures (degrees K) as a function of Fractional Cloud Amount (for cloud of 220 K and surface of 300 K) using $B(T) = (1-N)*B(T_{\text{sfc}}) + N*B(T_{\text{cld}})$.

Cloud Fraction N	Longwave Window		Shortwave Window		$T_s - T_1$
	Rad	Temp	Rad	Temp	
1.0	23.5	220	.005	220	0
.8	42.0	244	.114	267	23
.6	60.5	261	.223	280	19
.4	79.0	276	.332	289	13
.2	97.5	289	.441	295	6
.0	116.0	300	.550	300	0

SW and LW BTs for different cloud amounts
when $T_{cld}=220$ and $T_{sfc}=300$





Broken clouds appear different in 8.6, 11 and 12 um images;
 assume $T_{clr}=300$ and $T_{cld}=230$

$$T(11)-T(12)=[(1-N)*B_{11}(T_{clr})+N*B_{11}(T_{cld})]^{-1} \\ - [(1-N)*B_{12}(T_{clr})+N*B_{12}(T_{cld})]^{-1}$$

$$T(8.6)-T(11)=[(1-N)*B_{8.6}(T_{clr})+N*B_{8.6}(T_{cld})]^{-1} \\ - [(1-N)*B_{11}(T_{clr})+N*B_{11}(T_{cld})]^{-1}$$

Cold part of pixel has more influence at longer wavelengths

Band: 31 wavelength 11.00 μm



Instrument: MODIS

Lat = 39.655 Lon = 52.646

Band: 22 wavelength 3.97 μm



Instrument: MODIS

Lat = 39.655 Lon = 52.646

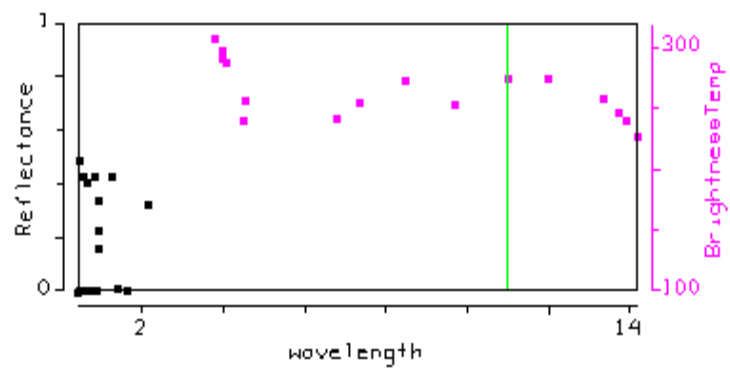
Cold clouds appear grainy in 4 μm MODIS images.

$$\Delta R = R_{\text{max}} / 2^{13} \quad \text{and} \quad \Delta T = \Delta R / [\text{dB}/\text{dT}]$$

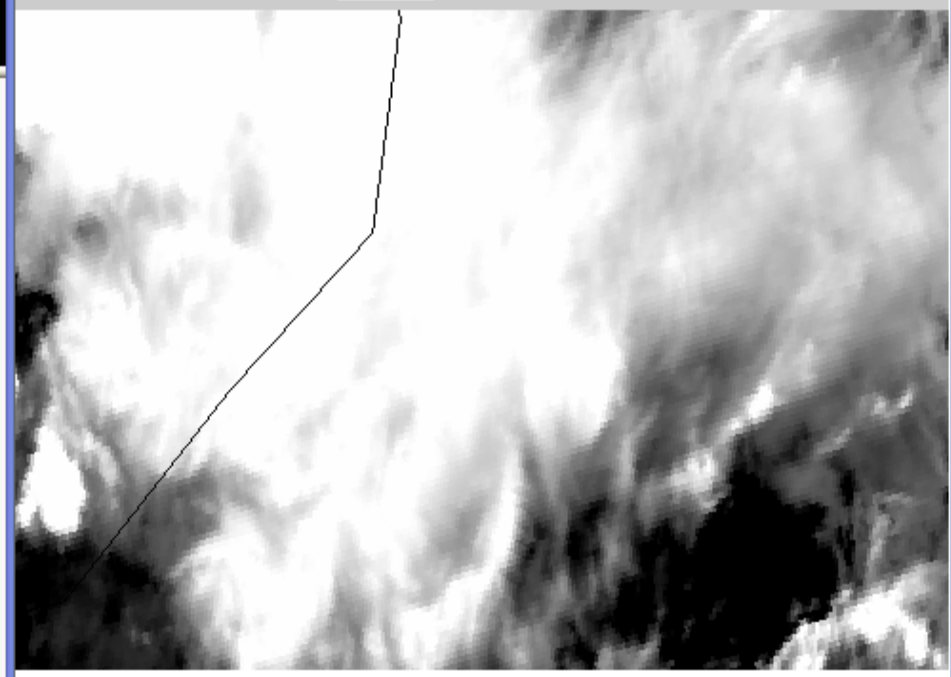
$\text{dB}/\text{dT}(4)$ is 100 times smaller at 200 K than at 300K;

Truncation error in cold scenes for 4 μm is several degrees K!

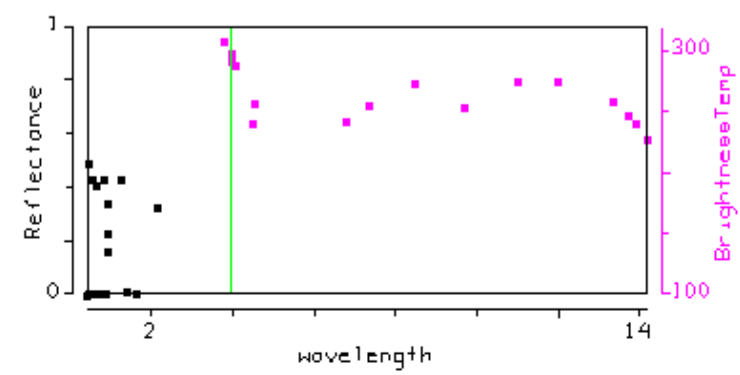
$\text{dB}/\text{dT}(11)$ is only 4 times smaller (hence it is not noticeable).



Band: 31 wavelength 11.00 μm



Instrument: MODIS Lat = 39.655 Lon = 52.646



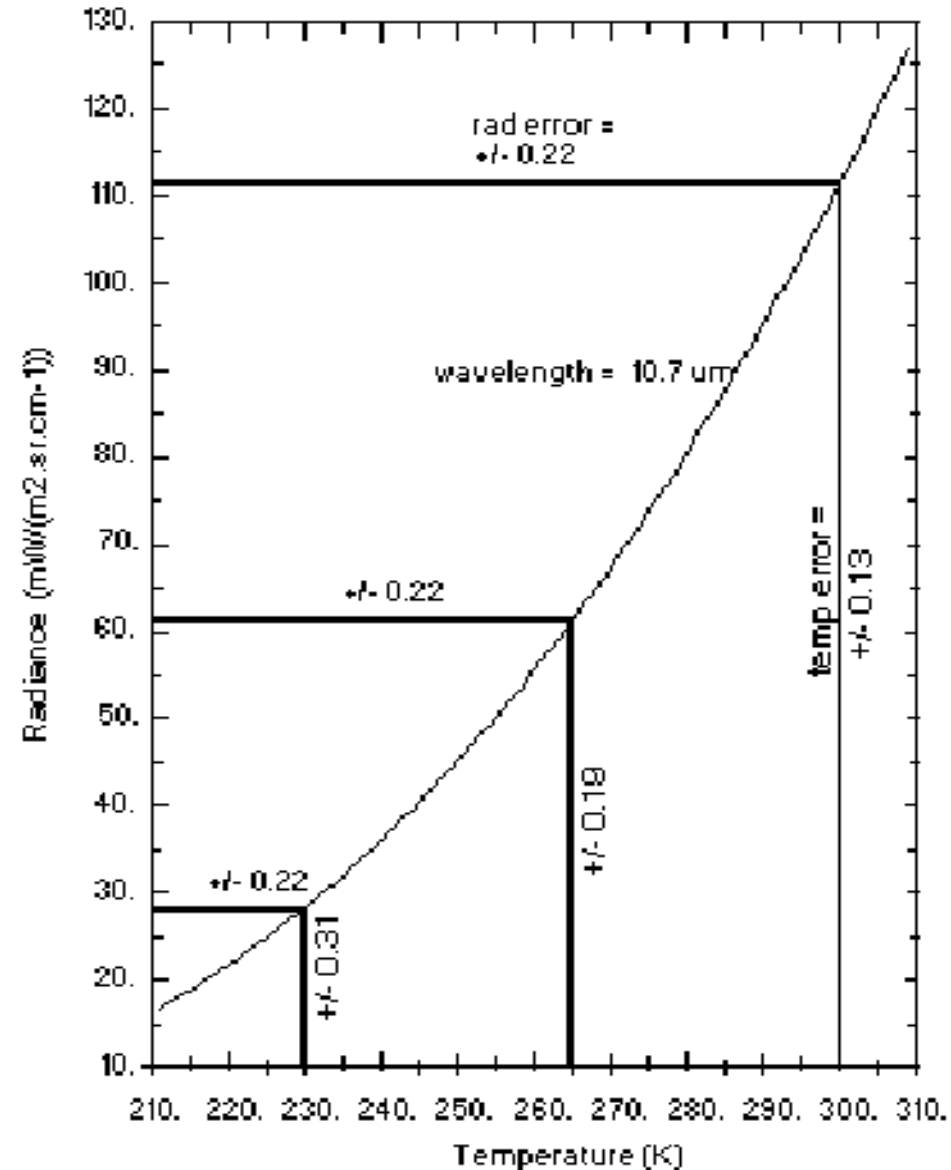
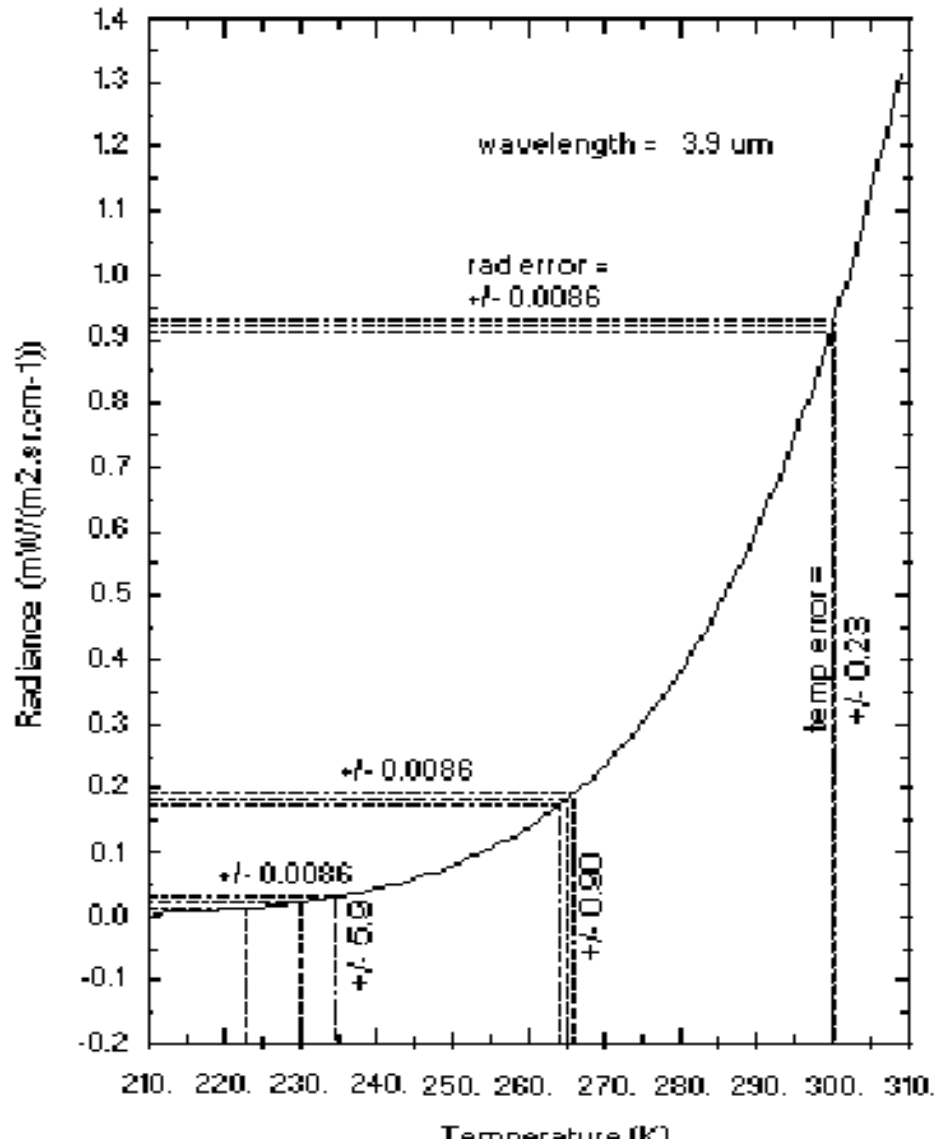
Band: 22 wavelength 3.97 μm



Instrument: MODIS Lat = 39.655 Lon = 52.646



NEDR vs NEDT at 4 and 11 μm



B and dB/dT at 4 and 11 μm

```
w,temp,r,drdt
rad = mW/ster/m2/cm-1
4      200      2.87612E-03  2.586404E-04
4      210      6.77262E-03  5.524178E-04
4      220      1.475345E-02  1.096472E-03
4      230      3.003489E-02  2.042302E-03
4      240      5.762783E-02  3.598814E-03
4      250      .1049535      6.040413E-03
4      260      .1825299      9.712637E-03
4      270      .3046972      1.503457E-02
4      280      .4903518      2.249792E-02
4      290      .7636536      3.266267E-02
4      300      1.154674      4.614974E-02
4      310      1.699957      6.363091E-02
4      320      2.442974      8.581714E-02
4      330      3.434435      .1134449
4      340      4.732509      .1472632
4      350      6.402821      .1880181
4      360      8.518434      .2364417
4      370      11.15956      .2932373
4      380      14.41332      .3590706
4      390      18.37315      .4345571
4      400      23.13851      .5202583
```

```
w,temp,r,drdt
rad = mW/ster/m2/cm-1
11     200     12.94316     .4238625
11     210     17.68172     .5254846
11     220     23.48358     .6363174
11     230     30.43456     .755097
11     240     38.60766     .8805164
11     250     48.06268     1.011279
11     260     58.84691     1.146147
11     270     70.99543     1.28396
11     280     84.53237     1.423662
11     290     99.4718      1.564306
11     300     115.8188     1.705055
11     310     133.5708     1.845184
11     320     152.7184     1.984071
11     330     173.2464     2.121195
11     340     195.1349     2.256122
11     350     218.3603     2.388497
11     360     242.8955     2.518038
11     370     268.7109     2.644525
11     380     295.7752     2.767794
11     390     324.0556     2.887726
11     400     353.5183     3.004245
```

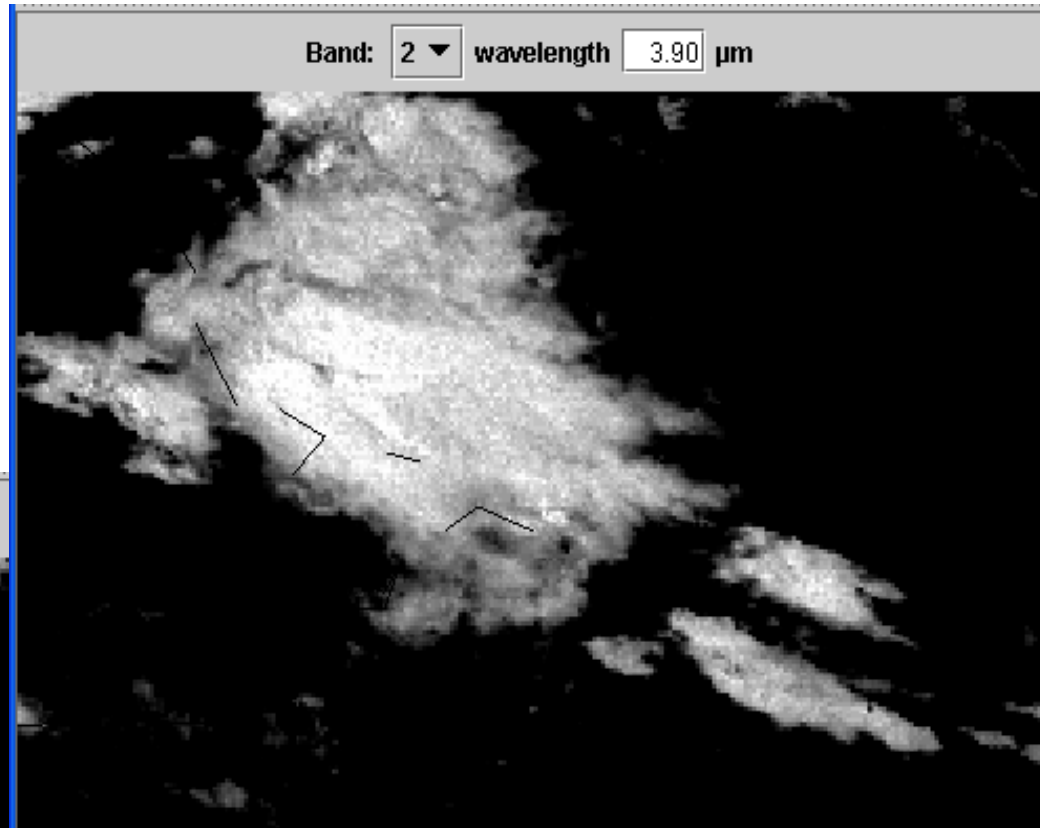
For GOES

with 10 bit data it is even worse

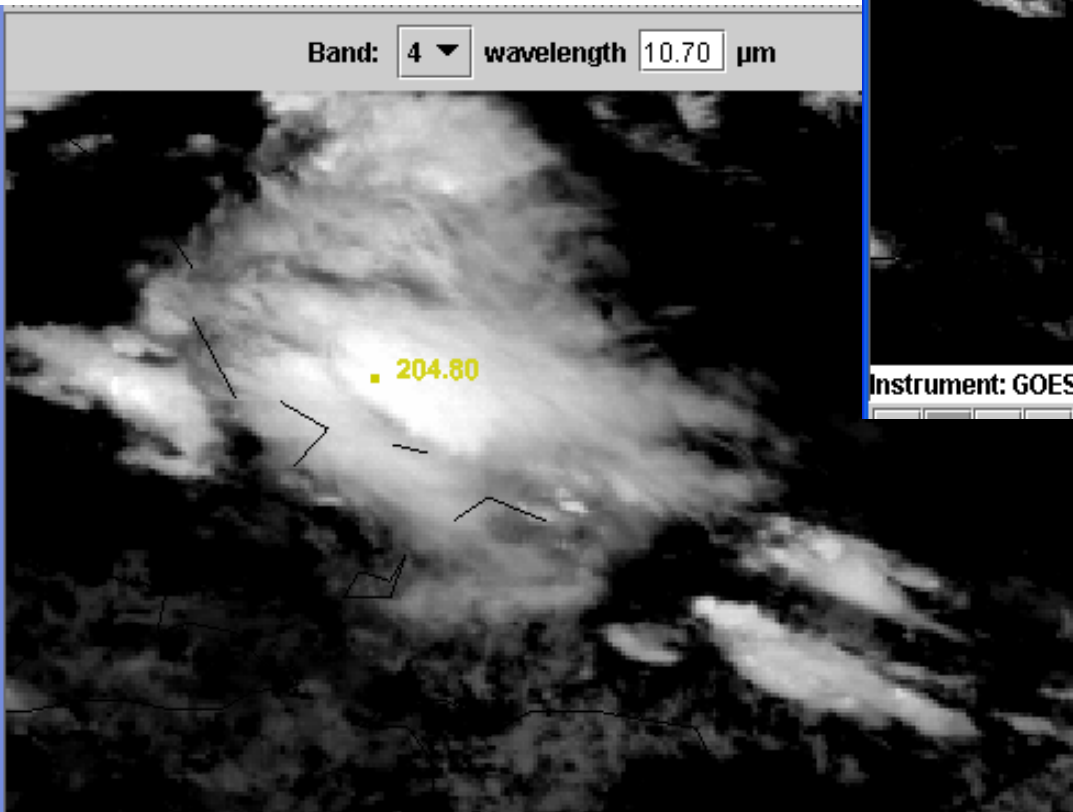
$$\Delta R = R_{\max} / 2^{10}$$

and

$$\Delta T = \Delta R / [\text{dB}/\text{dT}]$$



Instrument: GOES-E Lat = 13.931 Lon = -80.815



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